Research status of supercritical aviation kerosene and a convection heat transfer considering thermal pyrolysis

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Abstract

Purpose – This paper aims to comprehensively clarify the research status of thermal transport of supercritical aviation kerosene, with particular interests in the effect of cracking on heat transfer.

Design/methodology/approach – A brief review of current research on supercritical aviation kerosene is presented in views of the surrogate model of hydrocarbon fuels, chemical cracking mechanism of hydrocarbon fuels, thermo-physical properties of hydrocarbon fuels, turbulence models, flow characteristics and thermal performances, which indicates that more efforts need to be directed into these topics. Therefore, supercritical thermal transport of n-decane is then computationally investigated in the condition of thermal pyrolysis, while the ASPEN HYSYS gives the properties of n-decane and pyrolysis products. In addition, the one-step chemical cracking mechanism and SST k-ω turbulence model are applied with relatively high precision.

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Findings – The existing surrogate models of aviation kerosene are limited to a specific scope of application and their thermo-physical properties deviate from the experimental data. The turbulence models used to implement numerical simulation should be studied to further improve the prediction accuracy. The thermal-induced acceleration is driven by the drastic density change, which is caused by the production of small molecules. The wall temperature of the combustion chamber can be effectively reduced by this behavior, i.e. the phenomenon of heat transfer deterioration can be attenuated or suppressed by thermal pyrolysis.

Originality/value – The issues in numerical studies of supercritical aviation kerosene are clearly revealed, and the conjugation mechanism between thermal pyrolysis and convective heat transfer is initially presented.

Keywords Heat transfer deterioration, Supercritical aviation kerosene, Thermal pyrolysis, Thermal-induced acceleration

Paper type Research paper

Nomenclature

\( a \) = the size of a rectangular channel, mm;
\( A \) = the area of the cross-section of the channel, m\(^2\);
\( B_0^* \) = buoyancy factor;
\( c_p \) = specific heat of the fluid, J/kg·K;
\( c_s \) = volume concentration of component s, kmol/m\(^3\);
\( C_R \) = reactant concentration, kmol/m\(^3\);
\( d \) = hydraulic diameter of the circular pipe, m;
\( D_s \) = the diffusion coefficient of the component s, m\(^2\)/s;
\( D_{\text{wu}} \) = orthogonal divergence term;
\( E_a \) = activation energy and its value is 263.7 kJ/mol;
\( g \) = gravity acceleration, m/s\(^2\);
\( g_i \) = gravitational volume force along the \( i \) direction, m/s\(^2\);
\( G_{\text{k}} \) = term representing turbulent kinetic energy generated by laminar velocity gradient, m\(^2\)/s\(^2\);
\( G_{\omega} \) = term representing turbulent kinetic energy generated by \( \omega \) equation, m\(^2\)/s\(^2\);
\( Gr \) = Grashof number based on a temperature difference;
\( k \) = turbulent kinetic energy, m\(^2\)/s\(^2\);
\( k_0 \) = pre-exponential factor and its value is \( 1.6 \times 10^{15} \) 1/s;
\( K \) = reaction rate constant, 1/s;
\( K_v \) = acceleration parameter;
\( L \) = length of the horizontal rectangular channel, mm;
\( p \) = operating pressure, Pa;
\( \bar{p} \) = average pressure, Pa;
\( Pr_t \) = turbulent Prandtl number;
\( q_w \) = heat flux, kW/m\(^2\);
\( q_m \) = mass flow rate, kg/s;
\( R \) = universal gas constant and its value is \( 8.314 \times 10^{-3} \) kJ/mol·K;
\( Re \) = Reynolds number;
\( S_k \) = turbulent kinetic energy source term, m\(^2\)/s\(^2\);
\( S_s \) = the mass of the component produced per unit volume per unit time within the system by chemical reaction, i.e. production rate, kmol/m\(^3\)/s;
\( S_{\omega} \) = source term defined by users;
\( t \) = reaction time, s;
\( T \) = temperature, K;
\( \mathcal{T} \) = average temperature, K;
\( u \) = velocity of the fluid, m/s;
\( u_i \) = velocity of the fluid along with i direction, m/s;
\( \bar{u}_i \) = average velocity along with i direction, m/s;
\( \bar{u}_j \) = average velocity along with j direction, m/s;
\( v \) = inlet velocity, m/s;
\( x \) = x-axis;
\( x_i \) = x, y;
\( x_j \) = x, y;
\( y \) = y-axis;
\( Y_k \) = terms representing turbulence generated by diffusion; and
\( y^+ \) = dimensionless wall distance.

**Greek symbol**

\( \Gamma_k \) = diffusion rate of \( k \);
\( \Gamma_{\omega} \) = diffusion rate of \( \omega \);
\( \lambda \) = thermal conductivity of the fluid, W/m·K;
\( \lambda_t \) = turbulent thermal conductivity, W/m·K;
\( \mu \) = dynamic viscosity of the fluid, kg/m·s;
\( \mu_t \) = turbulent viscosity, kg/m·s;
\( \rho \) = density of the fluid, kg/m\(^3\);
\( \omega \) = specific dissipation; and
\( \omega_{K} \) = the reaction rate, kmol/m\(^3\)·s.

**Subscripts**

\( a \) = average;
\( b \) = bulk;
\( d \) = downstream section;
\( in \) = inlet;
\( out \) = outlet;
\( pc \) = pseudo-critical temperature;
\( t \) = tested section;
\( u \) = upstream section; and
\( w \) = wall.

**Acronyms**

AKN = Abe-Kondoh-Nagano;
DNS = direct numerical simulation;
HTD = heat transfer deterioration;
HTE = heat transfer enhancement;
LES = large eddy simulation;
LS = Launder–Sharma;
NIST = National Institute of Standards and Technology;
PPD = proportional product distribution;
RANS = Reynolds-averaged Navier–Stokes equations;
RP-3 = China no. 3 aviation kerosene;
1. Introduction

The scramjet engine can provide effective thrust for hypersonic vehicles (Sunden et al., 2016; Li et al., 2021a). However, thermal protection is an extremely crucial issue due to supersonic combustion and aerodynamic heating (Pizzarelli et al., 2016; Daniau et al., 2005). It is widely known that the active regenerative cooling technique (a schematic diagram is provided in Chen et al. (2020)) can be used to effectively solve this problem (Rao and Kunzru, 2006; Huang et al., 2002; Tsujikawa and Northam, 1996; Gascoin et al., 2008) and supercritical aviation kerosene is used in such a process. The heat transfer deterioration (HTD) and heat transfer enhancement (HTE) of supercritical hydrocarbon fuels can always be observed (Negoeescu et al., 2017; Xie et al., 2021; Kim and Kim, 2010); thus, the relevant heat transfer performance and flow characteristics were paid much attention to. Besides, based on the open published literature, the research in this area is summarized as shown in Figure 1(a), and it focuses on the mathematic model, impact of internal factors, impact of external factors as well as unconventional mechanism of supercritical fluid flow and heat transfer.

Regarding the supercritical hydrocarbon fuels, the numerical modelling and operating parameters are described with emphasis in this paper. The numerical models on supercritical hydrocarbon fuels include surrogate model, the chemical cracking mechanism, thermo-physical properties and turbulence models as shown in Figure 1(b).

1.1 Surrogate model of hydrocarbon fuels

Aviation kerosene is a compound with complex components, which cannot be thoroughly considered in the computational calculation. Thus, surrogate models of aviation kerosene are extensively used in the literature. Table 1 summarizes the surrogate models of aviation kerosene (Zhang et al., 2017; Fan and Yu, 2006; Ren et al., 2013; Zeng et al., 2014; Xiao et al., 2010; Zheng et al., 2015; Pei and Hou, 2017; Jiang et al., 2013; Cheng et al., 2012; Zhong et al., 2009a; Li et al., 2019a). Apart from literature Jiang et al. (2013), the rest are surrogate models of China no. 3 aviation kerosene (RP-3). Specifically, Zhang et al. (2017) investigated chemical recuperation in various aspect-ratio channels with n-decane as an alternative of aviation kerosene. In contrast, Fan and Yu (2006) put forward a three-component model, which was verified by an extended corresponding state model. The three-component model was composed of n-decane, trimethylcyclohexane and propylbenzene, which accounted for 0.49, 0.44 and 0.07, respectively, in the mole fraction. Ren et al. (2013) also reported a three-component surrogate model that has n-undecane (0.53 mole fraction), 1-butylcyclohexane (0.18 mole fraction) and 1,3,5-trimethyl-benzene (0.29 mole fraction), examined by the RP-3 experiment data. Zeng et al. (2014) proposed a three-component surrogate model including 0.65 volume fraction n-decane, 0.1 volume fraction toluene, 0.25 volume fraction propyl cyclohexane. Xiao et al. (2010) proposed a reaction model of kerosene using another three-component surrogate model, i.e. n-decane (0.79 mole fraction), trimethylcyclohexane (0.13 mole fraction) and ethylbenzene (0.08 mole fraction). Zheng et al. (2015) presented a four-component model, which can accurately evaluate the delayed time of ignition and the speed of laminar flame in real RP-3. The surrogate model includes n-decane, n-dodecane, ethycyclohexane and xylene, which have a share of 0.4, 0.42, 0.13 and 0.05, respectively in the mole fraction. Pei and Hou (2017) investigated the thermo-physical properties of a four-component surrogate model, which includes n-decane (0.35 mole fraction), n-undecane (0.20 mole fraction), 1-methyl-2-pentyl-cyclohexane (0.40 mole fraction) and 1,3,5-Trimethyl-benzene (0.05 mole fraction) in 300–1000 K and 1–15 MPa. Jiang et al. (2013) concluded that the HF-I kerosene can be replaced by a four-component model that
Figure 1.
(a) Summary of existing research on supercritical hydrocarbon fuel; (b) summary of numerical models of supercritical hydrocarbon fuel
Table 1. The surrogate models of aviation kerosene proposed by Chinese researchers.
<table>
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<tbody>
<tr>
<td>Benzenes</td>
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<tr>
<td>Toluene (C$_7$H$_8$)</td>
<td></td>
<td>–</td>
<td>–</td>
<td>0.1</td>
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<tr>
<td>Ethylbenzene (C$<em>8$H$</em>{10}$)</td>
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<td>–</td>
<td>0.08</td>
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<tr>
<td>p-Xylene (C$<em>8$H$</em>{10}$)</td>
<td></td>
<td>–</td>
<td>0.05</td>
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<tr>
<td>n-Propylbenzene (C$<em>9$H$</em>{12}$)</td>
<td></td>
<td>0.07</td>
<td>–</td>
<td>–</td>
<td>–</td>
<td>–</td>
<td>–</td>
<td>–</td>
<td>0.139</td>
<td>0.05</td>
<td>–</td>
<td>–</td>
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<tr>
<td>1,3,5-Trimethyl-benzene (C$<em>9$H$</em>{12}$)</td>
<td></td>
<td>0.29</td>
<td>–</td>
<td>–</td>
<td>–</td>
<td>–</td>
<td>0.05</td>
<td>–</td>
<td>–</td>
<td>–</td>
<td>0.09454</td>
<td>–</td>
</tr>
<tr>
<td>1-Methylnaphthalene (C$<em>{11}$H$</em>{10}$)</td>
<td></td>
<td>–</td>
<td>–</td>
<td>–</td>
<td>–</td>
<td>–</td>
<td>–</td>
<td>–</td>
<td>–</td>
<td>0.03</td>
<td>0.471</td>
<td>–</td>
</tr>
</tbody>
</table>
HFF

has n-decane (0.25 mole fraction), n-dodecane (0.50 mole fraction), n-tridecane (0.12 mole fraction) and n-butylcyclohexane (0.13 mole fraction). Additionally, many other surrogate models of RP-3 have been reported as shown in Table 1 (Cheng et al., 2012; Zhong et al., 2009a; Li et al., 2019a).

1.2 Chemical cracking mechanism of hydrocarbon fuels

It is widely known that the cracking reaction can be observed once the temperature is above 750 K (Edwards, 2006). The endothermic pyrolysis essentially affects supercritical thermal transport. Thus, it is critical to reveal the chemical cracking mechanism of hydrocarbon fuels. Based on the literature, the detailed reaction mechanism, lumped reaction mechanism and global reaction mechanism are commonly used as the cracking reaction models of supercritical hydrocarbon fuels. Furthermore, these cracking reaction models are embedded in commercial softwares such as ANSYS FLUENT, and this procedure is used to study the coupling of the cracking reaction, flow characteristics and heat transfer behavior. Ward et al. (2004, 2005) proposed two chemical cracking mechanisms of one-step proportional product distribution (PPD), and these models were used by Bao et al. (2014), Zhao et al. (2018), Tao et al. (2018a), Lei et al. (2020) and Sun et al. (2018a). Dahm et al. (2004) suggested 1175 reactions and 175 species of n-dodecane using the EXGAS software. Herbinet et al. (2007) also used the EXGAS software to report 1449 reactions and 271 species of n-dodecane. Xing et al. (2009) and Zhong et al. (2009b) studied the kinetic parameters of global reaction of RP-3 at \( p = 2.0 \text{–} 3.9 \text{ MPa, } T = 663 \text{–} 703 \text{ K and } p = 3.5 \text{–} 4.5 \text{ MPa, } T = 700 \text{–} 1100 \text{ K.} \) Jiang et al. (2013) presented a molecular reaction model with 28 species and 24 reactions, and the computational accuracy was acceptable when the cracking conversion was less than 86\%. Furthermore, this chemical cracking mechanism was adopted by Zhao et al. (2015), Zhang et al. (2015), Xu and Meng (2015a) and Jing et al. (2018) to study the impact of pressure on the fuel cracking, thermo-hydrodynamic characteristics and curved regenerative cooling channels, respectively. Hou et al. (2013) established a global reaction of one-step thermal cracking based on the proportional product distributions and gaseous-product cooling experiments. Ruan et al. (2014) developed a reduced 12-species chemical cracking mechanism, and it was also applied to analyze the effects of the endothermic fuel cracking on thermo-hydrodynamic characteristics (Sun et al., 2018b). Zhu et al. (2014) gave a global reaction framework including 18 species in conditions of \( p = 4.2 \text{–} 5.3 \text{ MPa, } T = 860 \text{–} 910 \text{ K.} \) Zhou et al. (2017) created a one-step global mild-cracking reaction framework for n-decane with 11 species. Wang et al. (2017) obtained a small-scale chemical reaction model of n-decane in conditions of 4 MPa, 480–720°C. Jiao et al. (2018) reported a molecular kinetic model with 15-steps reactions and only one primary reaction is included. Li et al. (2019b) proposed a simplified chemical cracking mechanism with 16 species and 11 reactions. The chemical cracking mechanisms mentioned above are summarized in Table 2.

1.3 Thermo-physical properties of hydrocarbon fuels

It is vital to obtain accurate properties of aviation kerosene for computational calculation. According to the existing literature, SUPERTRAPP, Aspen HYSYS, Aspen Plus and REFPROP, developed by the National Institute of Standards and Technology (NIST) (NIST, 1999), are always used to achieve the properties of aviation kerosene (Yang et al., 2020; Li et al., 2018; Bao et al., 2012; Jiao et al., 2020). The properties are also obtained by in-house codes based on extended corresponding state principles, which are always incorporated into the commercial softwares (Ruan et al., 2014; Tao et al., 2019; Wang et al., 2010). Besides, the Chemkin files (CHEMKIN-PRO 15092, 2009) can be imported into the ANSYS FLUENT to build the reaction model (Zhou et al., 2017). Herein, two things need to be confirmed. First,
<table>
<thead>
<tr>
<th>Authors (Year)</th>
<th>Reaction steps</th>
<th>Chemical cracking reactions</th>
<th>Measurement conditions</th>
<th>Activation energy $E_a$, kcal/mol (kJ/mol)</th>
<th>A factor $A$, s$^{-1}$ (% conversion)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Ward et al. (2004)</td>
<td>1-step (19-species)</td>
<td>C$<em>{10}$H$</em>{22}$ → 0.151 H$<em>2$ + 0.143 CH$<em>4$ + 0.256 C$<em>2$H$<em>4$ + 0.126 C$<em>3$H$<em>6$ + 0.230 C$<em>4$H$<em>6$ + 0.180 C$<em>5$H$<em>8$ + 0.196 C$<em>6$H$</em>{10}$ + 0.102 C$<em>7$H$</em>{12}$ + 0.171 C$<em>8$H$</em>{10}$ + 0.124 C$<em>9$H$</em>{12}$ + 0.195 C$</em>{10}$H$</em>{12}$ + 0.089 C$</em>{11}$H$</em>{14}$ + 0.169 C$</em>{12}$H$</em>{14}$ + 0.072 C$</em>{13}$H$</em>{16}$ + 0.152 C$</em>{14}$H$</em>{16}$ + 0.012 C$<em>{15}$H$</em>{18}$ + 0.053 C$<em>{16}$H$</em>{20}$ + 0.003 C$<em>{17}$H$</em>{20}$</td>
<td>$p = 3.45$ MPa, $T = 773$-873 K.</td>
<td>63</td>
<td>$(263.34)$</td>
</tr>
<tr>
<td>Adopted by others</td>
<td>1-step (18-species)</td>
<td>C$<em>{10}$H$</em>{22}$ → 0.153 CH$<em>4$ + 0.222 C$<em>2$H$<em>4$ + 0.138 C$<em>3$H$<em>6$ + 0.2 C$<em>4$H$<em>6$ + 0.185 C$<em>5$H$<em>8$ + 0.171 C$<em>6$H$</em>{10}$ + 0.118 C$<em>7$H$</em>{12}$ + 0.149 C$<em>8$H$</em>{10}$ + 0.137 C$<em>9$H$</em>{12}$ + 0.17 C$</em>{10}$H$</em>{12}$ + 0.091 C$</em>{11}$H$</em>{14}$ + 0.132 C$</em>{12}$H$</em>{16}$ + 0.046 C$</em>{13}$H$</em>{18}$ + 0.031 C$</em>{14}$H$_{20}$</td>
<td>$p = 3.45$ MPa, $T = 823$-873 K.</td>
<td>63</td>
<td>$(263.34)$</td>
</tr>
<tr>
<td>Adopted by others</td>
<td>1175-steps (175-species)</td>
<td>1175 reactions of n-dodecane using EXGAS software, it is difficult to couple detailed kinetic models with CFD code</td>
<td>$p = 0.1$ MPa, $T = 950$, 1000, 1050 K.</td>
<td>–</td>
<td>–</td>
</tr>
<tr>
<td>Adopted by others</td>
<td>None (271-species)</td>
<td>1449 reactions of n-dodecane using EXGAS software, it is difficult to couple detailed kinetic models with CFD code</td>
<td>$p = 0.1$ MPa, $T = 500$-1500 K.</td>
<td>–</td>
<td>–</td>
</tr>
<tr>
<td>Adopted by others</td>
<td>None (18-species)</td>
<td>Primary reaction: HF-I (C$<em>{11.85}$H$</em>{23.82}$) → $0.1086$H$_2$ + $0.4773$ CH$_4$ + $0.5586$C$_2$H$_4$ + $0.39$C$_3$H$_6$ + 0.41 C$_4$H$_8$ + 0.2001C$<em>5$H$</em>{10}$</td>
<td>$p = 5$ MPa, $T = 933$-973 K.</td>
<td>52.13</td>
<td>$(217.9)$</td>
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(continued)
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<th>Authors (Year)</th>
<th>Reaction steps</th>
<th>Chemical cracking reactions</th>
<th>Measurement conditions</th>
<th>Activation energy E&lt;sub&gt;a&lt;/sub&gt; kcal/mol (kJ/mol)</th>
<th>A factor A s&lt;sup&gt;-1&lt;/sup&gt;</th>
<th>(%) conversion</th>
</tr>
</thead>
<tbody>
<tr>
<td>Adopted by others</td>
<td>Hou et al. (2013)</td>
<td>1-step (9-species)</td>
<td>C&lt;sub&gt;12&lt;/sub&gt;H&lt;sub&gt;23&lt;/sub&gt; → 0.16 H&lt;sub&gt;2&lt;/sub&gt; + CH&lt;sub&gt;4&lt;/sub&gt; + 0.58 C&lt;sub&gt;2&lt;/sub&gt;H&lt;sub&gt;6&lt;/sub&gt; + 0.43 C&lt;sub&gt;3&lt;/sub&gt;H&lt;sub&gt;4&lt;/sub&gt; + 0.42 C&lt;sub&gt;3&lt;/sub&gt;H&lt;sub&gt;6&lt;/sub&gt; + 0.28 C&lt;sub&gt;4&lt;/sub&gt;H&lt;sub&gt;8&lt;/sub&gt; + 0.25 C&lt;sub&gt;4&lt;/sub&gt;H&lt;sub&gt;6&lt;/sub&gt; + 0.84 C&lt;sub&gt;4&lt;/sub&gt;H&lt;sub&gt;8&lt;/sub&gt;</td>
<td>p = 1–3 MPa, T = 780–988 K.</td>
<td>63 (263.34)</td>
<td>2.1 × 10&lt;sup&gt;15&lt;/sup&gt;</td>
</tr>
<tr>
<td>Adopted by others</td>
<td>Ruan et al. (2014)</td>
<td>1-step (12-species)</td>
<td>C&lt;sub&gt;10&lt;/sub&gt;H&lt;sub&gt;22&lt;/sub&gt; → 0.153 CH&lt;sub&gt;4&lt;/sub&gt; + 0.222 C&lt;sub&gt;2&lt;/sub&gt;H&lt;sub&gt;4&lt;/sub&gt; + 0.138 C&lt;sub&gt;2&lt;/sub&gt;H&lt;sub&gt;6&lt;/sub&gt; + 0.200 C&lt;sub&gt;3&lt;/sub&gt;H&lt;sub&gt;6&lt;/sub&gt; + 0.185 C&lt;sub&gt;3&lt;/sub&gt;H&lt;sub&gt;8&lt;/sub&gt; + 0.171 C&lt;sub&gt;4&lt;/sub&gt;H&lt;sub&gt;6&lt;/sub&gt; + 0.118 C&lt;sub&gt;5&lt;/sub&gt;H&lt;sub&gt;10&lt;/sub&gt; + 0.149 C&lt;sub&gt;5&lt;/sub&gt;H&lt;sub&gt;12&lt;/sub&gt; + 0.137 C&lt;sub&gt;6&lt;/sub&gt;H&lt;sub&gt;12&lt;/sub&gt; + 0.630 C&lt;sub&gt;6&lt;/sub&gt;H&lt;sub&gt;14&lt;/sub&gt; + 0.268 C&lt;sub&gt;7&lt;/sub&gt;H&lt;sub&gt;14&lt;/sub&gt;</td>
<td>–</td>
<td>63 (263.34)</td>
<td>1.6 × 10&lt;sup&gt;15&lt;/sup&gt;</td>
</tr>
<tr>
<td>Adopted by others</td>
<td>Sun et al. (2018b)</td>
<td>1-step (18-species)</td>
<td>C&lt;sub&gt;10&lt;/sub&gt;H&lt;sub&gt;22&lt;/sub&gt; → 0.0170 H&lt;sub&gt;2&lt;/sub&gt; + 0.1827 CH&lt;sub&gt;4&lt;/sub&gt; + 0.1980 C&lt;sub&gt;2&lt;/sub&gt;H&lt;sub&gt;6&lt;/sub&gt; + 0.0566 C&lt;sub&gt;3&lt;/sub&gt;H&lt;sub&gt;8&lt;/sub&gt; + 0.0372 C&lt;sub&gt;3&lt;/sub&gt;H&lt;sub&gt;6&lt;/sub&gt; + 0.0135 C&lt;sub&gt;3&lt;/sub&gt;H&lt;sub&gt;8&lt;/sub&gt; + 0.0048 C&lt;sub&gt;4&lt;/sub&gt;H&lt;sub&gt;10&lt;/sub&gt; + 0.2015 C&lt;sub&gt;5&lt;/sub&gt;H&lt;sub&gt;10&lt;/sub&gt; + 0.1167 C&lt;sub&gt;5&lt;/sub&gt;H&lt;sub&gt;12&lt;/sub&gt; + 0.3033 C&lt;sub&gt;6&lt;/sub&gt;H&lt;sub&gt;12&lt;/sub&gt; + 0.0735 C&lt;sub&gt;6&lt;/sub&gt;H&lt;sub&gt;14&lt;/sub&gt; + 0.2611 C&lt;sub&gt;6&lt;/sub&gt;H&lt;sub&gt;14&lt;/sub&gt; + 0.0684 C&lt;sub&gt;6&lt;/sub&gt;H&lt;sub&gt;16&lt;/sub&gt; + 0.2209 C&lt;sub&gt;6&lt;/sub&gt;H&lt;sub&gt;16&lt;/sub&gt; + 0.0053 C&lt;sub&gt;6&lt;/sub&gt;H&lt;sub&gt;16&lt;/sub&gt; + 0.0130 C&lt;sub&gt;6&lt;/sub&gt;H&lt;sub&gt;18&lt;/sub&gt; + 0.0817 C&lt;sub&gt;6&lt;/sub&gt;H&lt;sub&gt;18&lt;/sub&gt; + 0.0040 C&lt;sub&gt;6&lt;/sub&gt;H&lt;sub&gt;18&lt;/sub&gt;</td>
<td>p = 4.2–5.3 MPa, T = 860–910 K.</td>
<td>63 (263.34)</td>
<td>1.6 × 10&lt;sup&gt;15&lt;/sup&gt;</td>
</tr>
<tr>
<td>Adopted by others</td>
<td>Zhou et al. (2017)</td>
<td>None (11-species)</td>
<td>C&lt;sub&gt;10&lt;/sub&gt;H&lt;sub&gt;22&lt;/sub&gt; → 0.31 CH&lt;sub&gt;4&lt;/sub&gt; + 0.066 C&lt;sub&gt;2&lt;/sub&gt;H&lt;sub&gt;6&lt;/sub&gt; + 0.189 C&lt;sub&gt;2&lt;/sub&gt;H&lt;sub&gt;8&lt;/sub&gt; + 0.037 C&lt;sub&gt;3&lt;/sub&gt;H&lt;sub&gt;6&lt;/sub&gt; + 0.0135 C&lt;sub&gt;3&lt;/sub&gt;H&lt;sub&gt;6&lt;/sub&gt; + 0.0048 C&lt;sub&gt;4&lt;/sub&gt;H&lt;sub&gt;10&lt;/sub&gt; + 0.2015 C&lt;sub&gt;5&lt;/sub&gt;H&lt;sub&gt;10&lt;/sub&gt; + 0.1167 C&lt;sub&gt;5&lt;/sub&gt;H&lt;sub&gt;12&lt;/sub&gt; + 0.3033 C&lt;sub&gt;6&lt;/sub&gt;H&lt;sub&gt;12&lt;/sub&gt; + 0.0735 C&lt;sub&gt;6&lt;/sub&gt;H&lt;sub&gt;14&lt;/sub&gt; + 0.2611 C&lt;sub&gt;6&lt;/sub&gt;H&lt;sub&gt;14&lt;/sub&gt; + 0.0684 C&lt;sub&gt;6&lt;/sub&gt;H&lt;sub&gt;16&lt;/sub&gt; + 0.2209 C&lt;sub&gt;6&lt;/sub&gt;H&lt;sub&gt;16&lt;/sub&gt; + 0.0053 C&lt;sub&gt;6&lt;/sub&gt;H&lt;sub&gt;16&lt;/sub&gt; + 0.0130 C&lt;sub&gt;6&lt;/sub&gt;H&lt;sub&gt;18&lt;/sub&gt; + 0.0817 C&lt;sub&gt;6&lt;/sub&gt;H&lt;sub&gt;18&lt;/sub&gt; + 0.0040 C&lt;sub&gt;6&lt;/sub&gt;H&lt;sub&gt;18&lt;/sub&gt;</td>
<td>p = 3.5–4.0 MPa, T = 880–890 K.</td>
<td>64.3 (268.774)</td>
<td>1.75 × 10&lt;sup&gt;15&lt;/sup&gt;</td>
</tr>
<tr>
<td>Adopted by others</td>
<td>Wang et al. (2017)</td>
<td>None (16-species)</td>
<td>C&lt;sub&gt;10&lt;/sub&gt;H&lt;sub&gt;22&lt;/sub&gt; → 0.044 H&lt;sub&gt;2&lt;/sub&gt; + 0.186 CH&lt;sub&gt;4&lt;/sub&gt; + 0.241 C&lt;sub&gt;2&lt;/sub&gt;H&lt;sub&gt;6&lt;/sub&gt; + 0.0286 C&lt;sub&gt;3&lt;/sub&gt;H&lt;sub&gt;6&lt;/sub&gt; + 0.212 C&lt;sub&gt;4&lt;/sub&gt;H&lt;sub&gt;8&lt;/sub&gt; + 0.166 C&lt;sub&gt;4&lt;/sub&gt;H&lt;sub&gt;6&lt;/sub&gt; + 0.040 C&lt;sub&gt;5&lt;/sub&gt;H&lt;sub&gt;6&lt;/sub&gt; + 0.026 C&lt;sub&gt;5&lt;/sub&gt;H&lt;sub&gt;10&lt;/sub&gt; + 0.004 C&lt;sub&gt;6&lt;/sub&gt;H&lt;sub&gt;10&lt;/sub&gt; + 0.813 C&lt;sub&gt;5&lt;/sub&gt; + 0.001 C&lt;sub&gt;6&lt;/sub&gt;H&lt;sub&gt;16&lt;/sub&gt;</td>
<td>p = 4.0 MPa, T = 753–993 K.</td>
<td>59.4 (248.292)</td>
<td>6.209 × 10&lt;sup&gt;15&lt;/sup&gt;</td>
</tr>
</tbody>
</table>

(continued)
<table>
<thead>
<tr>
<th>Authors (Year)</th>
<th>Reaction steps</th>
<th>Chemical cracking reactions</th>
<th>Measurement conditions</th>
<th>Activation energy</th>
<th>A factor A, s⁻¹</th>
<th>(%) conversion</th>
</tr>
</thead>
<tbody>
<tr>
<td><strong>Adopted by others</strong></td>
<td>None</td>
<td>15-steps (12-species)</td>
<td>0.01047 H₂ + 0.01047 CH₄ + 0.01047 C₂H₆ + 0.01047 C₂H₄ + 0.01047 C₃H₈ + 0.01047 C₃H₆ + 0.01047 C₄H₁₀ + 0.01047 C₄H₈ + 0.01047 C₅ + 0.01047 CC₅ + 0.01047 C₆+</td>
<td>p = 3–5 MPa, T = 943.15–923.52 K.</td>
<td>58.54</td>
<td>3MPa: 4.52 × 10¹⁵, 4MPa: 5.22 × 10¹⁵, 5MPa: 5.69 × 10¹⁵</td>
</tr>
<tr>
<td><strong>Adopted by others</strong></td>
<td>None</td>
<td>11-steps (16-species)</td>
<td>0.07336 H₂ + 0.4710 CH₄ + 0.2981 C₂H₄ + 0.4174 C₂H₆ + 0.3287 C₃H₆ + 0.2483 C₄H₆ + 0.3007 C₅ + 0.1646 C₄H₁₀ + 0.0331 C₅H₆ + 0.6411 C₆ + 0.243 CC₅ + 0.05308 C₆H₂n+6</td>
<td>p = 3.5 MPa, T = 798–973 K.</td>
<td>56.6</td>
<td>1.223 × 10¹⁴</td>
</tr>
<tr>
<td><strong>Adopted by others</strong></td>
<td>None</td>
<td></td>
<td></td>
<td></td>
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</tbody>
</table>

Table 2. Supercritical aviation kerosene
the thermo-physical properties of the same substance calculated by different softwares must be consistent. Second, the thermo-physical properties of the different surrogate models of hydrocarbon fuels calculated by the same software must be consistent.

To start, n-decane is selected as the object of comparison. Its thermo-physical properties calculated by REFPROP and ASPEN HYSYS are compared and shown in Figure 2. It is obvious that the two softwares give similar density and dynamic viscosity. Besides, their specific heats are similar except for the maximum value. However, an obvious difference can be observed regarding the thermal conductivity. Thus, it should be careful to use the thermal properties of supercritical fluids, which need to be validated by the experimental data.

Then, the thermo-physical properties of the different surrogate models of hydrocarbon fuels calculated by the same software should be compared before the numerical simulations. Fortunately, this work has been partially done by Xu and Meng (2015b). The thermodynamic and transport properties of five surrogate models of RP-3 were examined including one-species, three-species in Fan et al. (2006), four-species in Mawid et al. (2003), six-species in Lenhert et al. (2007) and ten-species in Zhong et al. (2009a). It was found that the four-species model showed an excellent performance with experimental data, and it was regarded as the most preferable option. Surprisingly, the one-species model, saying n-decane, also showed an appropriate performance. Besides, Cheng et al. (2016) evaluated the three-species (Ren et al., 2013), three-species (Fan and Yu, 2006), four-species (Dagaut and Cathonnet, 2006), six-species (Violi et al., 2002), ten-species (Zhong et al., 2009a), and it was indicated that the surrogate model of three-species in Ren et al. (2013) can predict the thermo-physical properties well. In this work, the surrogate models of one-species in Zhang et al. (2017), three-species in Xiao et al. (2010), four-species in Jiang et al. (2013), five-species in Cheng et al. (2012) and ten-species in Zhong et al. (2009a) listed in Table 1 are carefully examined and the thermo-physical properties of n-decane along with the temperature under the conditions of $p = 3.0$ MPa.
physical properties of these models calculated by the ASPEN HYSYS are compared with the experimental data in Deng et al. (2011, 2012a, 2012b). This is because some important components cannot be found in NIST. As shown in Figures 3–7, for all the surrogate models of the hydrocarbon fuel, the calculated density and dynamic viscosity are consistent with the experimental data in Deng et al. (2011, 2012a, 2012b). However, for the specific heat, a large difference is found among different surrogate models, particularly in the supercritical-temperature region. Especially, the value of the supercritical temperature for the four-species surrogate model matches well with the experimental data due to the same position of the maximum specific heat. It is concluded that the usage of the properties of the hydrocarbon fuel should be carefully done and those data also should be validated by the experimental data as long as possible.

1.4 Turbulence models
The numerical methods such as Direct Numerical Simulation (DNS), Large Eddy Simulation (LES) and Reynolds-averaged Navier Stokes (RANS) with wall treatment have been assessed (Li et al., 2014; Bae and Yoo, 2005; Niceno and Sharabi, 2013; Deng et al., 2011, 2012a, 2012b). As shown in Figure 3, the comparison between the data from experiment (Deng et al., 2011, 2012a, 2012b) and data from ASPEN HYSYS for one-species surrogate model in Zhang et al. (2017)
Among them, the RANS is paid much attention because the current turbulence models in RANS are built based on constant physical properties at subcritical conditions (Li et al., 2014), so that they cannot correctly predict the supercritical heat transfer behavior due to the extremely dramatical thermo-physical properties. Therefore, the modifications of the turbulence models in RANS are intensively studied. Tao et al. (2018b) modified the LS (Launder–Sharma) low-Reynolds number turbulence model based on the density fluctuation. Li et al. (2016) improved the prediction accuracy of HTD by considering the density fluctuation and thermal expansion coefficient fluctuation in the AKN turbulence model (Abe et al., 1994). Jiang et al. (2018) considered the buoyancy-induced production of turbulent kinetic energy and the turbulent Prandtl number to improve the prediction ability of the HTD.
1.5 Flow characteristics and thermal performances

As the above-related issues are resolved, the flow characteristics and thermal performances of supercritical hydrocarbon fuels can be investigated. The experimental research and numerical research are described, respectively.

**Experimental research:** Zhong et al. (2009a) analytically derived heat transfer coefficients of a 10-species surrogate model in extensive temperature, pressure, and flow rate of 300–1000 K, 2.6–5.0 MPa, and 10–100 g/s, respectively. It was reported that the HTE phenomenon occurs once the supercritical condition appears. Zhang et al. (2012) studied thermal transport of hydrocarbon fuel in an upward-flowing tube, indicating that the HTD phenomenon appeared once the acceleration parameter ($K_v$) was smaller than $1.5 \times 10^{-8}$ or the buoyancy factor ($Bo^*$) was lower than $1.6 \times 10^{-10}$. Zhu et al. (2013) concentrated on the flow resistance of hydrocarbon fuel in a round tube placed horizontally with varied pressure, temperature and flow rates. It was concluded that high temperature and flow rates
corresponded to large pressure drops. Fu et al. (2017a, 2017b) analyzed the convective heat transfer in vertical and U-turn tubes with various pressure, heating conditions, flow rates and tube sizes, finding that the heat transfer intensity agreed well with the existing correlation when $T_b/T_{pc} < 0.80$ and the heat transfer intensity increased with the decreasing diameter of the bend. Huang et al. (2015) also studied the effect of the key parameters on the thermal transport of supercritical fuel. Wen et al. (2017a) addressed the effect of buoyancy force on supercritical thermal transport, which could be evaluated by the $Gr_{vl}/Gr_{th}$ (Petukhov et al., 1974). They also studied supercritical fuel in a vertical helical pipe, which presented that the heat transfer intensity of the outside was 31.5% larger than that of the inside due to the centrifugal secondary flow (Wen et al., 2017b). Jia et al. (2019) explored the effect of an initiator (nitropropane) on the hydrocarbon pyrolysis in a miniature pipe, and it was demonstrated that the pyrolysis temperature of n-decane was reduced by 100 K. Qin et al. (2013) clarified the parameter-dependent chemical recuperation process of supercritical hydrocarbon fuel. Liang et al. (2017) researched the impact of high pyrolysis on thermal

![Figure 6. Comparison between the data from experiment (Deng et al., 2011; Deng et al., 2012a; Deng et al., 2012b) and data from ASPEN HYSYS for five-species surrogate model in Cheng et al. (2012)](image-url)
performance and a critical heat flux of 0.4–0.5 MW/m² was confirmed. Liu et al. (2019a) innovatively proposed an online hybrid method to analyze the pyrolysis mechanism. Wang et al. (2019) introduced the formation mechanisms of thermo-acoustic instability in a vertical upward circular tube. Jiang et al. (2019) announced that an additional heat sink will be activated once thermal cracking happens. Liu et al. (2019b) compared the thermo-hydrodynamic characteristics at the conditions of asymmetric heating and uniform heating.

Most of the above studies focused on the flow characteristics and thermal performances of supercritical aviation kerosene as well as the cracking reactions. The influence mechanism of convective heat transfer with pyrolysis reaction is difficult to explain by experimental data. Thus, numerical simulations are required to use the impact of the cracking reaction on the thermal transport.

Numerical research: Many studies were numerically implemented to study the thermal transport of supercritical fuel but the cracking reaction was not fully considered (Li et al., 2019c; Li et al., 2020; Cheng et al., 2018; Pu et al., 2019). Fortunately, the cracking reaction was considered in some openly published literature. Hu et al. (2017) pointed out that the
pressure drop could be increased due to the thermal-induced acceleration caused by the cracking reaction. Gong et al. (2019) studied the secondary reaction in the thermal transport of supercritical fuel. It was evidenced that the HTD phenomenon could be reduced by the secondary exothermic reaction. Xu and Meng (2016), Xu et al. (2018b) examined the convective heat transfer as well as the cracking reaction and surface coking in a round pipe and a ribbed pipe for supercritical fuel. It was reported that heat transfer was improved by the extra heat sink because of the thermal cracking and thermal-induced acceleration. Also, the pyrolytic chemical reaction rate was increased and the pyrolytic surface coking was decreased due to the effect of the ribs.

Thermal pyrolysis has been considered in the thermal transport of supercritical fuel, as discussed above, but which mostly focused on operating conditions, e.g. pressure, inlet temperature, flow orientation, pipe size and heat/mass flux ratio. The real impacts of the thermal pyrolysis on thermo-hydrodynamic characteristics of supercritical fuel are, however, not completely revealed until now. Therefore, further studies should be executed to clarify the underlying mechanism. In this work, only one operating condition is chosen to detect the impact of thermal pyrolysis on thermal transport in a horizontal tube (the upper wall cooling duct of the combustion chamber) and the typical conditions are \( p = 3.0 \) MPa, \( T_{\text{in}} = 500 \) K, \( v_{\text{in}} = 0.7 \) m/s (\( q_m = 1.5 \) g/s), \( q_w = 1.5 \) MW/m\(^2\) (only one wall is heated), \( a = 2 \) mm (the size of a rectangular channel).

2. Computational domain
A detailed description of the numerical model has been illustrated in Li et al. (2021b), as also shown in Figure 8. Rectangular channels are normally used to wrap combustion chambers to cool the combustor wall. When a hypersonic vehicle is flying horizontally, the cooling channels near the upper and lower wall are different because of the influence of gravity. In this study, the cooling channel near the upper wall serves as the object of the study. Only one channel is considered in the convective heat transfer with thermal pyrolysis in terms of symmetry. Furthermore, a two-dimensional (2-D) computational domain is applied due to the huge computational complexity of thermal pyrolysis. The entire computational domain is divided into three parts: a 30 mm upstream section to guarantee that the turbulence is fully developed, a 100 mm test section to obtain the data of interest, a 30 mm downstream section to prevent the backflow. The inlet size of the cooling channel is selected as 2 mm. The horizontal flow is analyzed in the study and the gravity is strictly involved because of the importance of the buoyancy force (Wen et al., 2017a).

3. Numerical procedure
3.1 Governing equations
The fluid flow and heat transfer processes in a horizontal tube are governed by the mass/momentum/energy conservation equations. Also, the thermal pyrolysis is solved by the species mass conservation equation. These formulas are given below.

\[
\frac{\partial (\rho \overline{u}_i)}{\partial x_i} = 0
\]  

(1)

\[
\frac{\partial}{\partial x_j} (\rho \overline{u}_i \overline{u}_j) = -\frac{\partial \rho}{\partial x_i} + \frac{\partial}{\partial x_j} \left[ \mu \frac{\partial \overline{u}_i}{\partial x_j} - \rho \overline{u}_i \overline{u}_j \right] + \rho g_i
\]  

(2)
\[
\frac{\partial}{\partial x_i}(\rho u_i T) = \frac{\partial}{\partial x_i}\left[\left(\frac{\lambda}{c_p} + \frac{\mu_i}{Pr_i}\right) \frac{\partial T}{\partial x_i}\right]
\]

(3)

\[Pr_i = \frac{\mu_i c_p}{\lambda_i}\]

(4)

\[
\frac{\partial}{\partial x_i}(\rho u_i c_s) = \frac{\partial}{\partial x_i}\left(D_s \frac{\partial c_s}{\partial x_i}\right) + S_s
\]

(5)

Based on the open published literature (Zhao et al., 2017; Han et al., 2018) and validation described in Section 4, the SST k-\omega is requisitioned, and the corresponding formulas are supplied below:

\[
\frac{\partial(\rho k)}{\partial t} + \frac{\partial(\rho u_i k)}{\partial x_i} = \frac{\partial}{\partial x_i}\left[\left(\mu + \sigma_k u_i\right) \frac{\partial k}{\partial x_i}\right] + G_k - Y_k + S_k
\]

(6)

\[
\frac{\partial(\rho \omega)}{\partial t} + \frac{\partial(\rho u_i \omega)}{\partial x_i} = \frac{\partial}{\partial x_i}\left[\Gamma_{\omega} \frac{\partial \omega}{\partial x_i}\right] + G_\omega - Y_\omega + D_\omega + S_\omega
\]

(7)
3.2 Boundary conditions
The initial and boundary conditions are important for solving the governing equations, which are shown in Figure 9. Specifically, on the inlet of the channel, an inlet temperature and a mass flow rate are designated to be 500 K and 1.5 g/s ($v_{in} = 0.7$ m/s), respectively, while the operating pressure is 3.0 MPa. The boundary condition of outflow is used on the outlet of the channel. A heat flux of 1.5 MW/m² is applied to a single wall, while the other walls are set as adiabatic. Noting that the solid zone in this work is not considered. Only the flow characteristic and heat transfer in fluid zone are paid attention under the condition of single wall heating, i.e. nonuniform heating. This type of heating exists in the practical application.

3.3 Chemical kinetics model
As stated in Section 1.2, a large amount of chemical cracking mechanisms of hydrocarbon fuels are provided in previous studies. The PPD chemical cracking mechanism put forward by Ward et al. (2004) has been applied by many researchers. In this study, this chemical cracking model is also adopted and its expression is shown as follows:
The reaction rate of the thermal pyrolysis of supercritical n-decane can be written as:

$$\omega_R = \frac{dC_R}{dt} = -KC_R$$  \hspace{1cm} (9)

$$K = k_0 \cdot \exp\left(-\frac{E_a}{RT}\right)$$  \hspace{1cm} (10)

where $k_0$ is pre-exponential factor and its value is $1.6 \times 10^{15}$ 1/s, $E_a$ is activation energy and its value is 263.34 kJ/mol.

3.4 Thermo-physical properties

At present, n-decane is used as the surrogate model of aviation kerosene. The thermo-physical properties of reactant and pyrolysis products were calculated by the ASPEN HYSYS. Then, as for the mixture, the specific heat, the density and the thermal conductivity and viscosity were deduced by the volume-weighted mixing law, the mixing law and the mass-weighted mixing law, respectively. These mixing laws can be found in ANSYS FLUENT 2020 R1.

3.5 Numerical methods

Figure 10 shows the computational domain and a structured grid system. It should be noted that only the $y$-axis region from 70 to 72 mm and the $x$-axis region from 0 to 2 mm are plotted. To guarantee the accuracy of the simulation, enhanced wall treatment is employed in the region near the wall, and $y^+$ is less than unity. The governing equations are discretized and iteratively solved by the finite volume method. The second-order upwind difference scheme is used to spatially discretize the relevant variations. The pressure-velocity coupling is realized by the function of coupled. The stiff chemistry solver was used for the chemistry solver. The turbulence-chemistry interaction is solved by the eddy-dissipation concept. The spatial discretization of reactants and products was conducted by the first-order upwind scheme. FLUENT 2020 R1 is used, and the simulation is considered to be converged when the residuals reach $10^{-5}$.

4. Numerical assessment

4.1 Study of the turbulence models

The RANS model is used to analyze the thermo-hydrodynamic characteristics. Based on the open published literature (Zhao et al., 2017; Han et al., 2018), the low Reynolds number SST $k$-$\omega$ is suitable to simulate the supercritical thermal transport. To further validate the rationality and reliability of numerical simulations, the numerical results are compared with tested measurements (Liu et al., 2015) in which the n-decane flowing upward in a round tube was investigated. Figure 11 depicts the numerical results with SST $k$-$\omega$. Thermal pyrolysis was considered in one case while it was not considered in another case. Also, the inner wall temperature from the experiment was shown in Figure 11. As was expected, the numerical results with thermal pyrolysis are better than those without thermal pyrolysis. However, the HTD phenomenon was brought forward. From the
location of $x/d_{in} = 93.5$ to the $x/d_{in} = 366.3$ (the last measurement point along the tube), the maximum deviation between the experimental data and numerical data is 7.0\%, and this indicates that the SST $k$-$\omega$ is appropriate for the study of the supercritical thermal transport coupling with thermal cracking. Nevertheless, it needs to be declared that the cause of the deviation is closely relevant to the precision of thermo-physical properties and chemical cracking reactions.

4.2 Study of the mesh sensitivity
It was illustrated in Section 3.5 that the computational domain should be discretized by a structured mesh before the numerical simulation. Four sets of grids were drawn, and the outlet average temperature and velocity [these expressions are shown in equations (11) and (12)] were used to confirm the finest grid settings. As shown in Table 3, the mesh numbers of 15,900, 25,000, 39,000, 51,000 were detected, and it is reported that the mesh number of 39,000 was suitable to study the conjugation between thermal transport and cracking, and the deviations of the outlet average temperature and velocity are $-0.17$ and $0.32\%$, respectively, by comparing with the results in the case with a mesh number of 51,000.

$$T_b = \frac{\int_A \rho u c_p T dA}{\int_A \rho u c_p dA}$$

(11)
5. Results and discussion

5.1 The heated wall temperature

In the supercritical heat transfer, the wall temperature distribution is commonly used to determine if the HTD and HTE phenomena will happen (Edwards, 2003; Hobold and Silva, 2016). In this study, the wall temperature distributions of the case without thermal pyrolysis and the case with thermal pyrolysis are demonstrated in Figure 12. It is found that the wall temperature with thermal cracking is lower than that without thermal cracking. Furthermore, near the region of $y = 100\text{ mm}$, an HTD phenomenon appears when the thermal pyrolysis is not considered, while an HTE phenomenon occurs when the thermal cracking is involved. This implies that the cracking can effectively decrease the wall temperature of the combustor. In the initial stages of heat transfer, i.e. before the position of $y = 32\text{ mm}$, the wall temperatures of the two cases are the same because the thermal pyrolysis occurs when the temperature is larger than 770 K (Jia et al., 2014). It is clearly

$$v_{\text{ave}} = \frac{1}{A} \int_A u dA$$  \hspace{1cm} (12)
found that the wall temperature is smaller than 770 K before the position of $y = 32$ mm. Also, this phenomenon indirectly proves the accuracy of the numerical simulation.

### 5.2 Distribution of $y$-axis velocity

It is common sense that the thermal transport is relevant to the flow characteristics. Figure 13 compared the $y$-axis velocity distributions between the case with thermal pyrolysis and the case without thermal pyrolysis. It is found that the average $y$-axis velocity of the fluid is increased by the thermal pyrolysis particularly in the vicinity of the outlet region. This is an important reason corresponding to the reduced temperature under the participation of thermal pyrolysis as shown in Figure 12. Also, the thermal pyrolysis increases the gradient of the $y$-axis velocity along with the mainstream, implying that the acceleration is generated by the large gradient of the density, which enhances heat transfer, i.e. decreasing the wall temperature. There are two reasons for the change in density. The first one is the dramatic variation of supercritical hydrocarbon fuel and the second one is the generation of the new product due to the thermal pyrolysis. The dramatic variation of supercritical hydrocarbon fuel has been widely introduced and only the effects of reactant and product distribution on the density of the hydrocarbon fuel were studied in this work.

### 5.3 The effects of reactant and product distribution

Figure 14 gives the cracking rate of n-decane at the exit section along the $x$-axis. It can be seen that the average cracking rate is 13%, and it is less than 20% (PPD model proposed is valid within 20% as shown in Table 2). However, the maximum cracking rate at the exit section is 26% (the position near the heated wall) and it is slightly larger than 20%. In Zhang (2016), it is reported that the PPD model proposed by Ward et al. can be extended with cracking rate of 76%, thus it is still used in this work. Figure 15 plots the mass fractions of C$_{10}$H$_{22}$ (reactant) and C$_{6}$H$_{12}$ (product) for the case with thermal pyrolysis, respectively. It is found that the mass fraction of C$_{10}$H$_{22}$ decreases with the flow direction due to the thermal...
pyrolysis, especially, in the zone near the heated plane. Accordingly, the mass fraction of products such as \( \text{C}_6\text{H}_{12} \) increases along the mainstream in the vicinity of the heated plane. This means that more small molecules are produced and the small density of the mixture (reactant and products) is caused as shown in Figure 16. In comparison to the case without thermal pyrolysis, a larger variation of density can be found in the case with thermal pyrolysis. This implies that the dramatic variation of density caused by thermal pyrolysis is dominant in the thermal-induced acceleration. By comparing Figures 15 and 16, it is found that the change of the density coincides with the production of small molecules. This further confirms that the generation of acceleration is driven by the drastic density change because of the thermal pyrolysis.

![Figure 14. The cracking rate of n-decane at the exit section along the x-axis](image)

![Figure 13. The y-axis velocity distributions for (a) the case without thermal pyrolysis and (b) the case with thermal pyrolysis, respectively](image)
6. Conclusions

Supercritical fuel cooling is an excellent way to decrease the wall temperature of the combustor in a scramjet engine. In the cooling process, HTD and HTE can be observed. Furthermore, the thermal pyrolysis phenomenon can also be found with an increase of fluid temperature. Regarding these issues, a brief review of current research on supercritical aviation kerosene phenomena was organized first in views of the surrogate model of hydrocarbon fuels, chemical cracking mechanism of hydrocarbon fuels, thermo-physical properties of hydrocarbon fuels, turbulence model, flow characteristics and thermal performances. Then, numerical simulations were implemented to study the supercritical thermal transport of n-decane with thermal pyrolysis. The properties of the reactant and the pyrolysis products were calculated by the ASPEN HYSYS. The PPD chemical

Figure 15.
The mass fractions of C_{10}H_{22} (reactant) and C_{6}H_{12} (product) for the case with thermal pyrolysis

Figure 16.
The density distributions for (a) the case without thermal pyrolysis and (b) the case with thermal pyrolysis, respectively
cracking mechanism was applied. The SST k-ω was used because it has relatively high precisions. Based on the above contents, some findings are summarized as follows:

- The surrogate model and chemical cracking mechanism of hydrocarbon fuels are the key steps to explore the thermo-hydrodynamic characteristics of supercritical aviation kerosene. The existing models have corresponding application scope, which should be paid more attention to.

- The thermo-physical properties of existing surrogate models of aviation kerosene cannot match well with the experimental data. Also, quite different thermo-physical properties are obtained from SUPERTRAPP, Aspen HYSYS, Aspen Plus and REFPROP, and in-house code based on the extended corresponding state principle. This issue should be noticed.

- The modification of turbulence models is only based on the existing turbulence models and more work should be implemented to enhance the prediction accuracy of thermo-hydrodynamic characteristics of supercritical fluids.

- Numerous experimental studies were dedicated to studying the thermo-hydrodynamic characteristics of supercritical aviation kerosene and thermal pyrolysis. However, the coupling influence mechanism between the convective heat transfer and thermal pyrolysis is difficult to explain by the experimental data. Thus, it is necessary to numerically investigate the impact of thermal pyrolysis on thermal transport. Based on this study, the wall temperature of the combustion chamber can be effectively reduced by thermal pyrolysis, i.e. the HTD phenomenon can be attenuated or suppressed by thermal pyrolysis. This is because the thermal-induced acceleration is driven by the drastic density change, which is caused by the production of small molecules. This reason is more important than the properties of the reactant and products.

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